Adaptive finite difference solution for fractional boundary value problems

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Homogeneous Dirichlet BVP of the fractional Laplacian

Main task: We are concerned with the numerical solution of homogeneous Dirichlet boundary value problem (BVP):

$$egin{cases} (-\Delta)^s u = f, & ext{ in } \Omega \ u = 0, & ext{ in } \Omega^c \equiv \mathbb{R}^d \setminus \Omega \end{cases}$$

where Ω is a bounded domain of \mathbb{R}^d $(d \ge 1)$ and $(-\Delta)^s$ is the fractional Laplacian (operator) of order $s \in (0, 1)$.

Representations of the fractional Laplacian:

• Fourier transform representation

$$(-\Delta)^{s}u = \mathcal{F}^{-1}(|\boldsymbol{\xi}|^{2s}\mathcal{F}(u))$$

• Singular integral representation

$$(-\Delta)^{s}u(\mathbf{x}) = C_{d,s} \text{ p.v. } \int_{\mathbb{R}^{d}} \frac{u(\mathbf{x}) - u(\mathbf{y})}{|\mathbf{x} - \mathbf{y}|^{d+2s}} d\mathbf{y}$$
$$= C_{d,s} \text{ p.v. } \int_{\Omega} \frac{u(\mathbf{x}) - u(\mathbf{y})}{|\mathbf{x} - \mathbf{y}|^{d+2s}} d\mathbf{y} + C_{d,s}u(\mathbf{x}) \int_{\Omega^{c}} \frac{1}{|\mathbf{x} - \mathbf{y}|^{d+2s}} d\mathbf{y}$$

• Several other representations ...

Regularity of BVP's solution:

• Ros-Oton and Serra (2014):

$$u(\mathbf{x}) \sim dist(\mathbf{x}, \partial \Omega)^s$$
, near $\partial \Omega$
 $u(\mathbf{x}) \in C^\infty$, in interior

• Ros-Oton and Serra (2014):

$$\|u\|_{C^{s}(\mathbb{R}^{d})} \leq C \|f\|_{L^{\infty}(\Omega)}$$

• Acosta and Borthagaray (2017):

$$\begin{split} |u|_{H^{s+\frac{1}{2}-0}(\Omega)} \leq \begin{cases} C \|f\|_{C^{\frac{1}{2}-s}(\Omega)}, & 0 < s < \frac{1}{2} \\ C \|f\|_{L^{\infty}(\Omega)}, & s = \frac{1}{2} \\ C \|f\|_{C^{\beta}(\Omega)}, & \frac{1}{2} < s < 1 \\ \end{cases} \text{ for some } \beta > 0 \end{split}$$



(a) s = 0.5 (k = 0)

Motivations:

- Interesting in theory: 1st-order, 2nd-order, ... why not 0.5th-order or 1.5th-order?
- Statistical mechanics: the concentration of particles performing Brownian motion (random walks with short-range jumps) follows the standard diffusion equation while the concentration of particles performing Lévy flights (random walks with long-range jumps) satisfies a fractional diffusion equation.
- Quantum mechanics: Klein-Gordon operators $\sqrt{-\Delta + m^2}$ (tempered fractional Laplacian)
- It has been reported that fractional models can give more accurate description of underlying phenomena in image processing, finance, and biology, especially for anomalous dynamics (compared to dynamics satisfying Gaussian distribution).
- Good tool for use to model global interaction, long-range decay, and multiple scales.

References (Books, theses, and Reviews, incomplete list):

- I. Podlubny: Fractional Differential Equations, Academic Press, Inc., San Diego, CA, 1999.
- K. Diethelm: The Analysis of Fractional Differential Equations, Springer-Verlag, Berlin, 2010.
- M. K. Ishteva: Properties and Applications of the Caputo Fractional Operator. Master Thesis, Universität Karlsruke, 2005.
- Lischke et al., JCP (2020): "What is the fractional Laplacian? A comparative review with new results"

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- Main challenges in numerical solution:
 - High cost: the stiffness matrix is a full matrix
 - Slow convergence (less than or equal to 1st order), due to low regularity of the solution
- It has attracted considerable attention in the last decade from researchers and a variety of methods have been developed: Finite element, finite difference, spectral, meshfree, discrete Galerkin, Monte Carlo methods).

Finite difference methods (incomplete list)

- H. Wang and T. Basu (SIAM J. Sci. Comput., 2012): 1D, based on Grünwald-Letnikov formula
- Y. Huang and A. Oberman (SIAM J. Numer. Anal., 2014): 1D FDMs, prove $\mathcal{O}(h^{2-2s})$ in L_{∞} for smooth solutions and $\mathcal{O}(h^s)$ for non-smooth solutions
- \bullet S. Duo, H. van Wyk, and Y. Zhang (J. Comput. Phys., 2018): new 1D scheme and prove $\mathcal{O}(h^2)$ for smooth solutions
- N. Du, H. Sun, and H. Wang (Comput. Appl. Math., 2019): a volume penalized finite difference scheme and a preconditioned Krylov subspace iterative algorithm
- Z. Hao and R. Du (J. Comput. Phys., 2021): FDM and $\mathcal{O}(h^2)$ for smooth solutions in L_{∞} .
- R. Han and S. Wu (SIAM J. Numer. Anal., 2022): prove $O(|\log h|h^{2-s})$ for 0 < s < 1 and $O(h^{2s})$ for s < 2/3 in L_{∞}
- M. Chen, W. Deng, C. Min, J. Shi, and M. Stynes (2023): prove N^{-min(rs,2-2s)} on graded meshes (DOI: 10.13140/RG.2.2.10784.15361)

Finite element methods (incomplete list)

- G. Acosta and J. Borthagaray (SIAM J. Numer. Anal. 2017): linear finite elements, prove $\mathcal{O}(h^{\frac{1}{2}-0})$ for uniform meshes and $\mathcal{O}(N^{-\frac{1}{d}})$ (for 1/2 < s < 1) for graded meshes in H^s
- G. Acosta, F. Bersetche, and J. Borthagaray (Comput. Math. Appl. 2017): 2D FEM code
- G. Acosta, J. Borthagaray, and N. Heuer (IMA J. Numer. Anal. 2018): nonhomogeneous Dirichlet problems
- J. Borthagaray, L. Del Pezzo, and S. Martínez (J. Sci. Comput. 2018): O(h^{min(1,s+1/2)−0}) for quasi-uniform meshes and O(N^{-(1+s)/d)} for graded meshes in L². Eigenvalue problems.
- M. Ainsworth and C. Glusa (Comput. Methods Appl. Mech. Engrg. 2017, Contemporary computational mathematics 2018): a sparse approximation to the stiffness matrix and an efficient multigrid implementation. $\mathcal{O}(N^{-(1+s)/d})$ for adaptive meshes in L^2
- M. Faustmann, M. Karkulik, and J. Melenk (SIAM J. Numer. Anal. 2022): Local convergence for FEM

Numerical solution of homogeneous Dirichlet problems

Spectral, DG, collocation, and meshfree methods (incomplete list)

- G. Pang, W. Chen, and Z. Fu (J. Comput. Phys. 2015): RBF collocation (meshfree)
- X. Zhang, M. Gunzburger, L. Ju (Comput. Methods Appl. Mech. Engrg. 2016): Collocation method
- F. Song, C. Xu, and G. Karniadakis (SIAM J. Sci. Comput. 2017): spectral method for spectral fractional Laplacian
- Q. Du, L. Ju, and J. Lu (Math. Comp. 2019): DG for time dependent problems
- H. Antil, P. Dondl, and L. Striet (SIAM J. Sci. Comput. 2021): sinc function method for spectral fractional Laplacian
- J. Burkardt, Y. Wu, and Y. Zhang (SIAM J. Sci. Comput. 2021): meshfree pseudospectral method
- H. Li, R. Liu and L. Wang (Numer. Math. Theory Methods Appl. 2022): Hermite spectral-Galerkin method

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Weak formulation

$$\frac{C_{d,s}}{2}\int_{\mathbb{R}^d}\int_{\mathbb{R}^d}\frac{(u(\boldsymbol{x})-u(\boldsymbol{y}))(v(\boldsymbol{x})-v(\boldsymbol{y}))}{|\boldsymbol{x}-\boldsymbol{y}|^{d+2s}}d\boldsymbol{x}d\boldsymbol{y}=\int_{\Omega}fvd\boldsymbol{x},\quad\forall v\in\tilde{H}^s(\Omega)$$

or

$$\begin{aligned} \frac{\mathcal{C}_{d,s}}{2} \int_{\Omega} \int_{\Omega} \frac{(u(\boldsymbol{x}) - u(\boldsymbol{y}))(v(\boldsymbol{x}) - v(\boldsymbol{y}))}{|\boldsymbol{x} - \boldsymbol{y}|^{d+2s}} d\boldsymbol{x} d\boldsymbol{y} \\ &+ \mathcal{C}_{d,s} \int_{\Omega} \int_{\Omega^c} \frac{u(\boldsymbol{x})v(\boldsymbol{x})}{|\boldsymbol{x} - \boldsymbol{y}|^{d+2s}} d\boldsymbol{x} d\boldsymbol{y} = \int_{\Omega} f v d\boldsymbol{x}, \quad \forall v \in \tilde{H}^s(\Omega) \end{aligned}$$

Features of finite element approximation:

- Need to compute 2d-dimensional singular integrals and integrals on unbounded domains.
- Full stiffness matrix: costly to assemble
- Full stiffness matrix: costly to carry out its multiplication with vectors: $\mathcal{O}(N_v^2)$ flops.
 - 2D: for a mesh of size $10^2 \times 10^2$, $N_v = 10^4$ and $N_v^2 = 10^8$
 - 3D: for a mesh of size $10^2 \times 10^2 \times 10^2$, $N_\nu = 10^6$ and $N_\nu^2 = 10^{12}$
- Works for arbitrary bounded domains and with mesh adaptation
- Sharp error estimates available: $\mathcal{O}(h^{\min(1,s+\frac{1}{2})-0})$ for quasi-uniform meshes and $\mathcal{O}(\bar{h}^{1+s})$ for graded and adaptive meshes in L^2 , $\bar{h} = N^{-1/d}$

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Fourier transform representation:

$$(-\Delta)^{s} u = \mathcal{F}^{-1}(|\boldsymbol{\xi}|^{2s}\mathcal{F}(u))$$

In 2D:

$$(-\Delta)^{s}u(x,y) = \frac{1}{(2\pi)^{2}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} (\widehat{-\Delta})^{s}u(\xi,\eta) e^{ix\xi} e^{iy\eta} d\xi d\eta$$
$$\widehat{(-\Delta)^{s}}u(\xi,\eta) = (\xi^{2} + \eta^{2})^{s} \hat{u}(\xi,\eta)$$

Discrete Fourier transform (DFT): Consider a uniform infinite grid (lattice)

$$(x_j, y_k) = (jh_{\text{FD}}, kh_{\text{FD}}), j, k \in \mathbb{Z},$$

DFT and its inverse on this grid are given by

$$\begin{split} \check{u}(\xi,\eta) &= \sum_{j=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} u_{j,k} e^{-ix_j \xi} e^{-iy_k \eta} \\ u(x_j,y_k) &= \frac{h_{\text{FD}}^2}{(2\pi)^2} \int_{-\frac{\pi}{h_{\text{FD}}}}^{\frac{\pi}{h_{\text{FD}}}} \int_{-\frac{\pi}{h_{\text{FD}}}}^{\frac{\pi}{h_{\text{FD}}}} \check{u}(\xi,\eta) e^{ix_j \xi} e^{iy_k \eta} d\xi d\eta \\ &= \frac{1}{(2\pi)^2} \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} \check{u}(\frac{\xi}{h_{\text{FD}}},\frac{\eta}{h_{\text{FD}}}) e^{ij\xi} e^{ik\eta} d\xi d\eta \end{split}$$

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Uniform-grid FD approximation based on Fourier transform

Consider the 5-point FD approximation to the Laplacian:

$$(-\Delta_h)u(x_j, y_k) = \frac{1}{h_{\text{FD}}^2}(u_{j+1,k} - 2u_{j,k} + u_{j-1,k}) + \frac{1}{h_{\text{FD}}^2}(u_{j,k+1} - 2u_{j,k} + u_{j,k-1})$$

Applying the DFT to the above equation, we get

$$(-\check{\Delta}_h)u(\xi,\eta) = \frac{1}{h_{\text{FD}}^2} \left(4\sin^2(\frac{\xi h_{\text{FD}}}{2}) + 4\sin^2(\frac{\eta h_{\text{FD}}}{2})\right) \check{u}(\xi,\eta).$$

The FD approximation of the fractional Laplacian is given by

$$(-\Delta_{h})^{s}u(x_{j},y_{k}) = \frac{h_{\text{FD}}^{2}}{(2\pi)^{2}} \int_{-\frac{\pi}{h_{\text{FD}}}}^{\frac{\pi}{h_{\text{FD}}}} \int_{-\frac{\pi}{h_{\text{FD}}}}^{\frac{\pi}{h_{\text{FD}}}} \frac{1}{h_{\text{FD}}^{2s}} \left(4\sin^{2}(\frac{\xi h_{\text{FD}}}{2}) + 4\sin^{2}(\frac{\eta h_{\text{FD}}}{2})\right)^{s} \check{u}(\xi,\eta)e^{ij\xi}e^{ik\eta}d\xi d\eta$$

where

$$\check{u}(\xi,\eta) = \sum_{j=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} u_{j,k} e^{-ix_j\xi} e^{-iy_k\eta}$$

Uniform-grid FD approximation based on Fourier transform

The FD approximation of the fractional Laplacian:

$$(-\Delta_{h})^{s}u(x_{j}, y_{k}) = \frac{1}{h_{\text{FD}}^{2s}} \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} A_{(j,k),(m,n)}u_{m,n} = \frac{1}{h_{\text{FD}}^{2s}} \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} T_{j-m,k-n}u_{m,n}$$
$$= \frac{1}{h_{\text{FD}}^{2s}} \sum_{m=-N}^{N} \sum_{n=-N}^{N} T_{j-m,k-n}u_{m,n} \quad (\text{ for } u = 0 \text{ on } \mathbb{R}^{d} \setminus \Omega)$$

• $T_{p,q}$'s are the Fourier coefficients of $\left(4\sin^2(\frac{\xi}{2}) + 4\sin^2(\frac{\eta}{2})\right)^s$, i.e.,

$$T_{p,q} = \frac{1}{(2\pi)^2} \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} \left(4\sin^2(\frac{\xi}{2}) + 4\sin^2(\frac{\eta}{2}) \right)^s e^{ip\xi} e^{iq\eta} d\xi d\eta$$

$$T_{-p,-q} = T_{p,q}, \quad T_{-p,q} = T_{p,q}, \quad T_{p,-q} = T_{p,q}$$

They can be approximated using FFT (fast Fourier transform)

• A_{FD} is block Teoplitz matrix, symmetric and positive definite, and

$$\lambda_{\min}(A_{\text{FD}}) \geq Ch_{\text{FD}}^{2s}, \quad \lambda_{\max}(A_{\text{FD}}) \leq C$$

- $A_{\text{FD}}u$ can be computed using FFT in $\mathcal{O}(N^d \log N^d)$ flops (almost linear about N^d)
- Works only for rectangular/cubic domains; cannot be incorporated with mesh adaptation.

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GoFD (J. Shen and WH: 2024) for arbitrary bounded domains

Existing methods:

- Finite difference methods: Efficient (using FFT), uniform rectangular meshes, simple domains
- Finite element methods: Slow (with full stiffness matrix), unstructured meshes, arbitrary domains

Aim: develop a method: efficient (via FFT), unstructured meshes, arbitrary domains

$$(-\Delta)^{s} u = f \text{ in } \Omega, \quad u = 0 \text{ in } \Omega^{c}$$
$$\implies \quad \frac{1}{h_{\text{FD}}^{2s}} D_{h}^{-1} (I_{h}^{\text{FD}})^{T} A_{\text{FD}} I_{h}^{\text{FD}} u = f \qquad (1)$$

where

- A_{FD} : uniform-grid FD approx. of $(-\Delta)^s$ on \mathcal{T}_{FD}
- I_h^{FD} : transfer matrix from \mathcal{T}_h to \mathcal{T}_{FD} (sparse)
- D_h : a diagonal matrix formed by the column sums of $I_h^{\rm FD}$



- $u = \{u(x_j), j = 1, ..., N_v\}, f = \{f(x_j), j = 1, ..., N_v\}, x_j$'s: vertices of \mathcal{T}_h
- (1) can be written into a symmetric system as

$$(I_h^{\text{FD}})^T A_{\text{FD}} I_h^{\text{FD}} \boldsymbol{u} = h_{\text{FD}}^{2s} D_h \boldsymbol{f}$$
⁽²⁾

Theorem 1 (J. Shen and WH, 2024)

If I_{h}^{FD} has full column rank, then D_{h} is invertible, $(I_{h}^{FD})^{T}A_{FD}I_{h}^{FD}$ is symmetric and positive definite, and thus (2) is solvable.

Theorem 2 (J. Shen and WH, 2024)

Let I_h^{FD} be the transfer matrix associated with piecewise linear interpolation. If the uniform grid's spacing h_{FD} satisfies

$$h_{FD} \leq rac{a_{min}}{(d+1)\sqrt{d}},$$

where a_{min} is the minimum height of T_h , then I_h^{FD} has full column rank. In this case, $(I_h^{FD})^T A_{FD} I_h^{FD}$ is symmetric and positive definite and thus invertible.

Features of GoFD:

• $(I_h^{\text{FD}})^T A_{\text{FD}} I_h^{\text{FD}}$ is symmetric and positive definite and satisfies

$$\lambda_{\min}((I_h^{\mathsf{FD}})^{\mathsf{T}} A_{\mathsf{FD}} I_h^{\mathsf{FD}}) \geq C h_{\mathsf{FD}}^{2s}(a_{\min}/h)^{2d}, \quad \lambda_{\max}((I_h^{\mathsf{FD}})^{\mathsf{T}} A_{\mathsf{FD}} I_h^{\mathsf{FD}}) \leq C$$

• $(I_h^{\text{FD}})^T A_{\text{FD}} I_h^{\text{FD}} \boldsymbol{u}$ can be carried out via FFT and therefore,

$$(I_h^{\mathsf{FD}})^T A_{\mathsf{FD}} I_h^{\mathsf{FD}} \boldsymbol{u} = h_{\mathsf{FD}}^{2s} D_h \boldsymbol{f}$$

can be solved efficiently using Krylov subspace iteration methods (e.g. Conjugate Gradient).

- Preconditioners based on ILU (J. Shen and WH 2024).
- Works for any bounded domain and with mesh adaptation.
- Convergence order will be demonstrated by numerical examples.

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Mesh terminology:

- $\mathcal{T}_h = \{K\}$: a simplicial mesh for $\Omega \subset \mathbb{R}^d$ (d = 1, 2, 3, ...)
- N: the number of elements, N_v : the number of vertices, \hat{K} : the master element
- $F_{\mathcal{K}}: \hat{\mathcal{K}} \to \mathcal{K}$ is the affine mapping and $F'_{\mathcal{K}}$ is the Jacobian matrix of $F_{\mathcal{K}}$.



- Edge matrix of K: $E_K = [\mathbf{x}_1^K \mathbf{x}_0^K, ..., \mathbf{x}_d^K \mathbf{x}_0^K]$
- Edge matrix of $\hat{\mathcal{K}}$: $\hat{\mathcal{E}} = [\boldsymbol{\xi}_1 \boldsymbol{\xi}_0, ..., \boldsymbol{\xi}_d \boldsymbol{\xi}_0]$
- Relation between F'_K , E_K , and \hat{E} :

$$\mathbf{x} = F_{\mathcal{K}}(\boldsymbol{\xi}) = F_{\mathcal{K}}'(\boldsymbol{\xi} - \boldsymbol{\xi}_0^{\mathcal{K}}) + \mathbf{x}_0^{\mathcal{K}} \implies E_{\mathcal{K}} = F_{\mathcal{K}}'\hat{E} \implies F_{\mathcal{K}}' = E_{\mathcal{K}}\hat{E}^{-1}$$

- $\mathbb{J}_{K} = (F'_{K})^{-1} = \hat{E}E_{K}^{-1}$
- Use the metric tensor (monitor function) $\mathbb{M} = \mathbb{M}(x)$ to control mesh concentration on Ω .

The MMPDE moving mesh method

Given a meshing energy/functional $I_h[\mathcal{T}_h] = \sum_{K} |K| G(\mathbb{J}_K, \det(\mathbb{J}_K), \mathbb{M}_K)$, the MMPDE approach defines the mesh nodal velocities as a gradient system of the energy:

$$\frac{d\boldsymbol{x}_{i}}{dt} = -\frac{P_{i}}{\tau} \left(\frac{\partial I_{h}}{\partial \boldsymbol{x}_{i}}\right)^{T} = \frac{P_{i}}{\tau} \sum_{K \in \omega_{i}} |K| \boldsymbol{v}_{i_{K}}^{K}, \quad i = 1, ..., N_{v}$$

• au > 0 is used to control the response time of the mesh movement to the changes in $\mathbb M$

- $P_i > 0$ is chosen to make the equation invariant under the scaling transformation of \mathbb{M}
- ω_i is the element patch associated with x_i
- i_K is the local index of x_i on K
- $\mathbf{v}_{i\kappa}^{K}$'s are the local velocities given by

$$\begin{bmatrix} \left(\mathbf{v}_{1}^{K}\right)^{T} \\ \vdots \\ \left(\mathbf{v}_{d}^{K}\right)^{T} \end{bmatrix} = -GE_{K}^{-1} + E_{K}^{-1}\frac{\partial G}{\partial \mathbb{J}}\hat{E}E_{K}^{-1} + \frac{\partial G}{\partial \det(\mathbb{J})}\frac{\det(\hat{E})}{\det(E_{K})}E_{K}^{-1} - \frac{1}{d+1}\sum_{j=0}^{d}\operatorname{tr}\left(\frac{\partial G}{\partial \mathbb{M}}\mathbb{M}_{j,K}\right) \begin{bmatrix} (E_{K}^{-1})_{j} \\ \vdots \\ (E_{K}^{-1})_{j} \end{bmatrix}$$

$$\left(\mathbf{v}_{0}^{K}\right)^{T} = -\sum_{k=1}^{d}\left(\mathbf{v}_{k}^{K}\right)^{T} - \sum_{j=0}^{d}\operatorname{tr}\left(\frac{\partial G}{\partial \mathbb{M}}\mathbb{M}_{j,K}\right) (E_{K}^{-1})_{j}$$

$$\mathbb{M}_{j,K} = \mathbb{M}(\mathbf{x}_{j}^{K}), \qquad (E_{K}^{-1})_{j} : \text{ the } j\text{-th row of } E_{K}^{-1}$$

Theorem 6.1 (Kamenski & WH, Math Comput 2018)

Assume that the metric tensor is bounded

 $\underline{m} \mathbb{I} \leq \mathbb{M}(\mathbf{x}) \leq \overline{m} \mathbb{I}, \quad \forall \mathbf{x} \in \Omega$

for some positive constants \underline{m} and \overline{m} , and the meshing functional satisfies the coercivity condition

$$M_h = \sum_{\mathcal{K}} |\mathcal{K}| \mathcal{G}(\mathbb{J}_{\mathcal{K}}, \mathsf{det}(\mathbb{J}_{\mathcal{K}}), \mathbb{M}_{\mathcal{K}}): \qquad \mathcal{G}(\mathbb{J}, \mathsf{det}(\mathbb{J}), \mathbb{M}, \mathbf{x}) \geq lpha \left[\mathsf{tr}(\mathbb{J}\mathbb{M}^{-1}\mathbb{J}^{\mathcal{T}}) \right]^q - eta, \quad \forall \mathbf{x} \in \Omega,$$

with q > d/2 and positive constants α and β . If the initial mesh is nonsingular, then

- The mesh governed by the x-formulation of MMPDE will be nonsingular for t > 0;
- Specifically, the minimal height and volume of K are bounded below by

$$\begin{aligned} a_{K} &\geq C(\mathcal{T}_{h}^{0}) \ \overline{m}^{-\frac{d}{2(2q-d)} - \frac{1}{2}} N^{-\frac{2q}{d(2q-d)}}, \quad \forall K \in \mathcal{T}_{h}, \quad \forall t > 0 \\ |K| &\geq C(\mathcal{T}_{h}^{0}) \ \overline{m}^{-\frac{d^{2}}{2(2q-d)} - \frac{d}{2}} N^{-\frac{2q}{d(2q-d)}}, \quad \forall K \in \mathcal{T}_{h}, \quad \forall t > 0 \end{aligned}$$

- **9** Holds for fully discrete MMPDE if Δt is sufficiently small (depending on \overline{m} and N).
- e Holds for any (convex or concave) domain in any dimension.
- Works for the equidistribution-alignment functional and other meshing functionals.

M-uniform mesh approach and equidistribution & alignment:

- \mathcal{T}_h is said to be \mathbb{M} -uniform if it is uniform in the metric $\mathbb{M} = \mathbb{M}(\mathbf{x})$.
- Equivalently, all of its elements have the same size and are similar to \hat{K} , measured in \mathbb{M} .
- An M-uniform mesh satisfies the equidistribution condition (for element size)

$$|\mathcal{K}|\sqrt{\det(\mathbb{M}_{\mathcal{K}})} = rac{1}{N}\sigma_h, \quad orall \mathcal{K} \in \mathcal{T}_h \qquad \left(\ \sigma_h = \sum_{\mathcal{K}} |\mathcal{K}|\sqrt{\det(\mathbb{M}_{\mathcal{K}})} \
ight)$$

and the alignment condition (for element similarity - shape & orientation)

$$\frac{1}{d}\operatorname{tr}((F'_{K})^{T}\mathbb{M}_{K}F'_{K}) = \operatorname{det}((F'_{K})^{T}\mathbb{M}_{K}F'_{K})^{\frac{1}{d}}, \quad \forall K \in \mathcal{T}_{h}$$

• An energy functional associated with these conditions is $(\mathbb{J}_{\mathcal{K}} = (F'_{\mathcal{K}})^{-1}, \mathbb{M}_{\mathcal{K}} = \mathbb{M}(\mathbf{x}_{\mathcal{K}}))$

$$\begin{split} \mathcal{I}_{h}[\mathcal{T}_{h}] &= \sum_{K} |\mathcal{K}| \sqrt{\det(\mathbb{M}_{K})} \left(\theta \left(\operatorname{tr}(\mathbb{J}_{K} \mathbb{M}_{K}^{-1} \mathbb{J}_{K}^{T}) \right)^{\frac{dp}{2}} + (1 - 2\theta) d^{\frac{dp}{2}} \left(\frac{\det(\mathbb{J}_{K})}{\sqrt{\det(\mathbb{M}_{K})}} \right)^{p} \right) \\ &\equiv \sum_{K} |\mathcal{K}| \mathcal{G}(\mathbb{J}_{K}, \det(\mathbb{J}_{K}), \mathbb{M}_{K}) \end{split}$$

where $\theta \in (0, 0.5]$ and p > 1 are dimensionless parameters. Magic choice: $\theta = 1/3$, p = 2.

The MMPDE moving mesh method

The derivatives of G wrt \mathbb{J} , det(\mathbb{J}), and \mathbb{M} can be found using scalar-by-matrix differentiation as

$$\begin{split} G &= \theta \sqrt{\det(\mathbb{M})} \left(\operatorname{tr}(\mathbb{J}\mathbb{M}^{-1}\mathbb{J}^{T}) \right)^{\frac{dp}{2}} + (1-2\theta)d^{\frac{dp}{2}} \sqrt{\det(\mathbb{M})} \left(\frac{\det(\mathbb{J})}{\sqrt{\det(\mathbb{M})}} \right)^{\rho} \\ &\frac{\partial G}{\partial \mathbb{J}} = dp\theta \sqrt{\det(\mathbb{M})} \left(\operatorname{tr}(\mathbb{J}\mathbb{M}^{-1}\mathbb{J}^{T}) \right)^{\frac{dp}{2}-1} \mathbb{M}^{-1}\mathbb{J}^{T} \\ &\frac{\partial G}{\partial \det(\mathbb{J})} = p(1-2\theta)d^{\frac{dp}{2}} \det(\mathbb{M})^{\frac{1-\rho}{2}} \det(\mathbb{J})^{p-1} \\ &\frac{\partial G}{\partial \mathbb{M}} = -\frac{dp\theta}{2} \sqrt{\det(\mathbb{M})} \left(\operatorname{tr}(\mathbb{J}\mathbb{M}^{-1}\mathbb{J}^{T}) \right)^{\frac{dp}{2}-1} \mathbb{M}^{-1}\mathbb{J}^{T}\mathbb{J}\mathbb{M}^{-1} \\ &+ \frac{\theta}{2} \sqrt{\det(\mathbb{M})} \left(\operatorname{tr}(\mathbb{J}\mathbb{M}^{-1}\mathbb{J}^{T}) \right)^{\frac{dp}{2}} \mathbb{M}^{-1} \\ &+ \frac{(1-\rho)(1-2\theta)d^{\frac{dp}{2}}}{2} \det(\mathbb{M})^{\frac{1-\rho}{2}} \det(\mathbb{J})^{p}\mathbb{M}^{-1} \\ &P_{i} = \det(\mathbb{M}(\mathbf{x}_{i}))^{\frac{p-1}{2}} \end{split}$$

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Example 1

Example 1:

$$\begin{cases} (-\Delta)^{s} u = \frac{2^{2s} \Gamma(1+s+k) \Gamma(\frac{d}{2}+s+k)}{k! \Gamma(\frac{d}{2}+k)} \cdot P_{k}^{s,\frac{d}{2}-1} (2|\mathbf{x}|^{2}-1), & \text{in } \Omega \equiv B_{1}(0) \\ u = 0, & \text{in } B_{1}(0)^{c} \end{cases}$$

where $B_1(0)$ is the unit ball and $P_k^{s,\frac{d}{2}-1}(\cdot)$ is the Jacobi polynomial of degree k with parameters $(s, \frac{d}{2} - 1)$. The exact solution is

$$u = (1 - |\mathbf{x}|^2)^{\alpha}_+ P_k^{\alpha, \frac{d}{2} - 1} (2|\mathbf{x}|^2 - 1)$$

Remarks

- Adaptive meshes are obtained with the Moving Mesh PDE method (MMPDE) (Russell, Ren, and WH 1994 and 2010).
- Results are shown for k = 0 and k = 5 in 1D, 2D, and 3D

Example 1 - 1D

For quasi-uniform meshes: $\mathcal{O}(h^s)$ in L^{∞} norm and $\mathcal{O}(h^{\min(1,0.5+s)})$ in L^2 norm, consistent with convergence order proved and numerically demonstrated for finite element approximation



For adaptive meshes: $\mathcal{O}(\bar{h}^{s+0.5})$ in L^{∞} norm and $\mathcal{O}(\bar{h}^2)$ in L^2 norm, better than convergence order proved for finite element approximation. $\bar{h} = 1/N$.





Error in L^2 norm: $\mathcal{O}(h^{\min(1,0.5+s)})$ for quasi-uniform meshes and $\mathcal{O}(\bar{h}^2)$ for adaptive meshes



Adaptive meshes



Error in L^2 norm: $\mathcal{O}(h^{\min(1,0.5+s)})$ for quasi-uniform meshes and $\mathcal{O}(\bar{h}^{1+s})$ for adaptive meshes



Example 2 - 2D (Batman-shaped domain with s = 0.75)

$$\left\{ egin{array}{ll} (-\Delta)^s u = 1, & ext{in } \Omega \ u = 0, & ext{in } \Omega^c \end{array}
ight.$$

Analytical solution is not available. The reference solution (for computing the error) is obtained with a fine adaptive mesh. Complex domain with holes inside and waving exterior boundary.



The error in L^2 norm:

- $\mathcal{O}(h^{\min(1,0.5+s)})$ for quasi-uniform meshes
- $\mathcal{O}(\bar{h}^2)$ for adaptive meshes

Example 3 - 2D (L-shaped domain with s = 0.5)

$$\left\{ egin{array}{ll} (-\Delta)^s u = 1, & ext{in } \Omega \ u = 0, & ext{in } \Omega^c \end{array}
ight.$$

Analytical solution is not available. The reference solution (for computing the error) is obtained with a fine adaptive mesh.



The error in L^2 norm:

- $\mathcal{O}(h^{\min(1,0.5+s)})$ for quasi-uniform meshes
- $\mathcal{O}(\bar{h}^2)$ for adaptive meshes

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Conclusions

GoFD for Dirichlet problems of the fractional Laplacian:

- Relatively simple to code
- Efficient implementation via FFT
- Works for arbitrary bounded domains
- Able to incorporate with mesh adaptation strategies
- ullet Solvability guaranteed when $\mathit{h}_{\mathsf{FD}} \leq \mathit{a}_{\min}/(d+1)/\sqrt{d}$
- Effective preconditioners based on ILU and circulant matrix
- Numerical examples show: GoFD has convergence order $\mathcal{O}(h^{\min(1,0.5+s)})$ for quasi-uniform meshes and $\mathcal{O}(\bar{h}^2)$ for adaptive meshes, both in L^2 norm.

References

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Thank you for your attention!